## Molecular Beam Epitaxy Grown Group-IV Alloys: Ideal Candidate for Momentum(k)-Space Carrier Separation Photodetectors

<u>Tyler T. McCarthy<sup>1\*</sup></u>, Zheng Ju<sup>1,2</sup>, Stephen Schaefer<sup>1</sup>, Xin Qi<sup>1</sup>, Allison McMinn<sup>1</sup>, Shui-Qing Yu<sup>3</sup>, and Yong-Hang Zhang<sup>1</sup>

<sup>1</sup>School of Electrical, Computer, and Energy Engineering, Arizona State University, Tempe, AZ 85287, USA

<sup>2</sup>Department of Physics, Arizona State University, Tempe, AZ 85287, USA

<sup>3</sup>Department of Electrical Engineering, University of Arkansas, Fayetteville, Arkansas 72701, USA \*Corresponding author: <u>ttmccart@asu.edu</u>

Recently, we proposed the momentum(k)-space carrier separation (k-SCS) concept that combines the advantages of both direct and indirect bandgaps for light detection/conversion devices. The basic principle is to have a direct bandgap in a semiconductor that has a slightly larger bandgap than the indirect fundamental bandgap, giving a thermalization barrier,  $\Delta_{\Gamma-L}$ , such as  $3k_BT$ , for electrons at the conduction band edge. In this example, the  $\Gamma$ -valley minimum is  $3k_BT$  higher in energy than the L-valley minimum. The sharp absorption edge of the direct bandgap energy,  $E_{g,\Gamma}$ , appears at just slightly (3 $k_BT$ ) higher in energy than the slow onset of the absorption edge at the indirect bandgap,  $E_{g,L}$ . Under light illumination, electrons are excited to the direct  $\Gamma$ -valley in the conduction band while leaving holes in the valence band edge at the symmetry point of  $\Gamma$ . The large majority of photogenerated electrons in the direct  $\Gamma$ -valley will quickly thermalize at a sub-picosecond time scale to the lower energy indirect L-valley. These electrons in the L-valley will recombine with the holes in the valence band at a time scale of tens of microseconds to milliseconds. Both carriers, electrons and holes, are transported in real space but with different momentums, i.e. separately in k-space, to their corresponding contacts with negligible recombination. This clever design not only improves photogenerated carrier lifetime, similar to indirect bandgap semiconductors, but also offers a large absorption coefficient, similar to direct bandgap semiconductors.

Group-IV alloys such as GeSn are a model material system to demonstrate the novel idea as photodetectors with Sn compositions near the indirect-to-direct bandgap transition are predicted to have greater detectivity than conventional IV-VI and III-V compound photodetectors at room temperature, and comparable detectivity to InAs detectors operating at 77 K. GeSn samples with  $\Delta_{\Gamma-L}$  between  $0.4k_BT$  and  $3k_BT$  were grown by MBE on Ge substrates for MWIR (2 to 5 µm). Substrate surfaces were first cleaned using HF and HCl solution prior to a UHV outgas at 550 °C. A Ge buffer was grown at a substrate temperature of 500 °C before cooling down to 200 °C for GeSn growth. Ge effusion cell was held constant while Sn effusion cell was varied between 825 to 900 °C to obtain designed composition. RHEED showed a streaky (2x1) surface reconstruction pattern that transitions to a (1x1) with increasing Sn. Introducing Si expands the wavelength coverage range to 2 ~ 22 µm, making it ideal for MWIR, LWIR and VLWIR applications. More details of the theory and experiments will be reported at the conference.



Fig. 1. Schematic illustrating momentum(*k*)-space carrier separation using  $Si_xGe_{1-x-y}Sn_y$  as an example material. a) An electron thermalization barrier,  $\Delta_{\Gamma-L}$ , on the order of  $k_BT$  is introduced between the  $\Gamma$ - and *L*-valleys in the conduction band. The barrier height  $\Delta_{\Gamma-L}$  is given by the difference between the indirect bandgap (*L*- $\Gamma$  transition,  $E_{g,L}$ ) and the direct bandgap ( $\Gamma$ - $\Gamma$  transition,  $E_{g,\Gamma}$ ). Photogenerated electrons are excited mainly into the  $\Gamma$  valley due to higher absorption coefficient of the direct  $\Gamma$ -  $\Gamma$  transition (G), and then quickly thermally relax to the indirect *L* valley at a lower energy, where recombination (R) takes place at a much lower rate, resulting in a long lifetime. b) The absorption coefficient of the higher energy direct bandgap is much greater than the indirect bandgap which lies several  $k_BT$  below.



Fig. 2. Calculated bandgap energies ( $\Gamma$ - $\Gamma$ , L- $\Gamma$ , X- $\Gamma$ ) of the SiGe, GeSn, and SiSn binaries as a function of their lattice constants at 300 K. The shaded region indicates the SiGeSn alloys with direct bandgaps that satisfy the requirement  $0.4k_BT \le \Delta_{\Gamma-L} \le 3k_BT$ , a range for the thermalization barrier.



Fig. 3. (a,b) RHEED reconstruction changes from (2x) to (1x) with increasing Sn composition and (c) XRD spectra of GeSn samples towards achieving *k*-SCS. (inset) Sn composition versus Sn cell temperature calibration curve for constant substrate and Ge cell temperatures